

# Quantum State and Process Tomography with Entangled Photons

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# Optical Approaches to Quantum Information

The polarization of a photon comprises an ideal 2-state quantum system, equivalent to any other 2-state quantum system (e.g., two-level atom, spin-1/2, etc.), but often much easier to produce, manipulate, and detect.

- Quantum communication – q. cryptography, dense coding, teleportation
- Quantum “network” – coupling together remote quantum computers, multiparty protocols
- Small-scale quantum circuits – algorithms, repeaters, DFS studies, simulating other systems
- *Scalable* quantum computing?

Current goals:

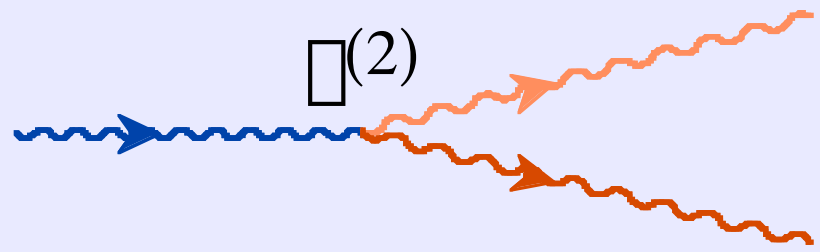
- push the limits on control and measurement (error  $< 10^{-4}$ )
- develop a “*benchmark technology*”
- explore fundamentals of quantum info. processing.

Basic tools: arbitrary quantum state synthesis, precise process control, accurate q. tomography

Advanced tools: quantum storage, single-photon sources, entanglement-on-demand, etc.

# Spontaneous Parametric Downconversion

Burnham & Weinberg, PRL **25**, 84 (1970)



The diagram shows a blue wavy line with an arrow pointing right, labeled  $\chi^{(2)}$  above it. This line splits into two orange wavy lines with arrows pointing right and slightly away from each other.

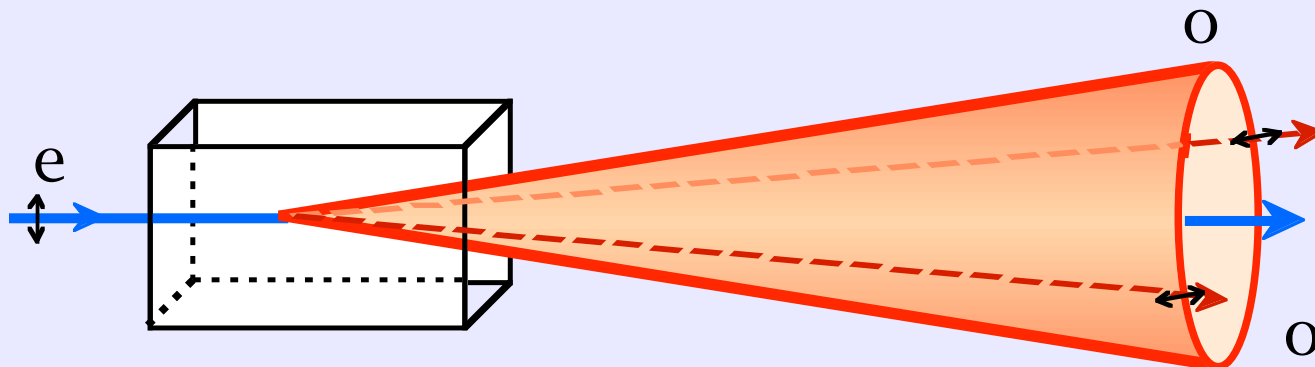
$$\chi_p = \chi_s + \chi_i^*$$
$$\bar{\chi}_p = \bar{\chi}_s + \bar{\chi}_i^\dagger$$

\*Energy conservation  $\rightarrow$  energy entanglement

$\dagger$  Momentum conservation  $\rightarrow$  momentum entanglement

## Type-I Phase-matching

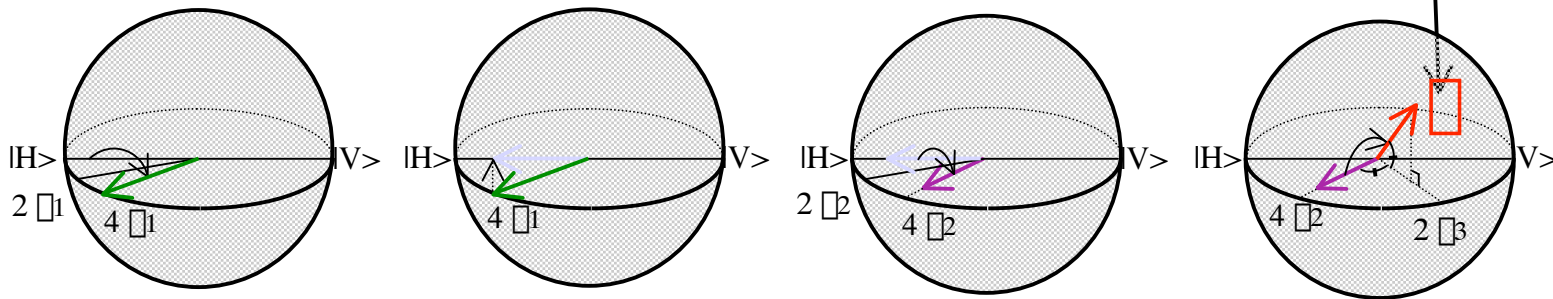
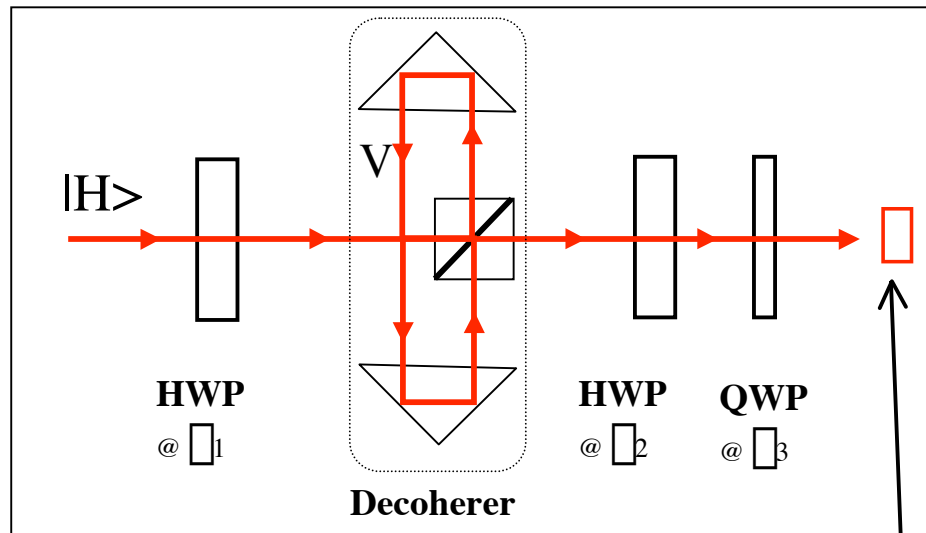
(Photons have identical polarizations)



# Single Qubit Generation

An arbitrary state:

$$|\psi\rangle = \begin{bmatrix} A \\ B e^{i\phi} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & e^{i\phi} \end{bmatrix} \begin{bmatrix} A \\ B \end{bmatrix}$$



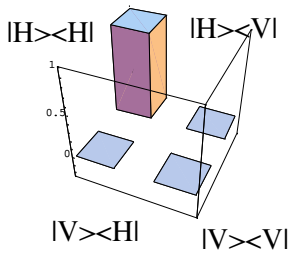
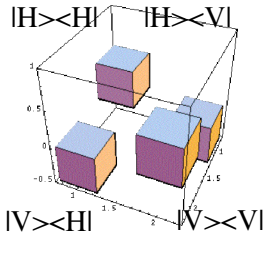
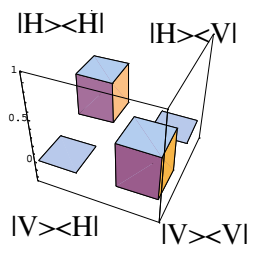
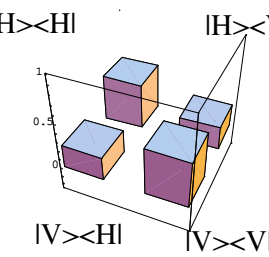
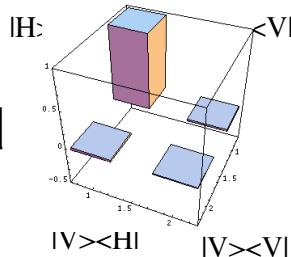
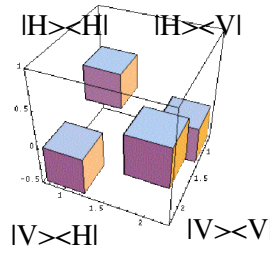
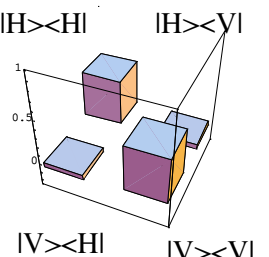
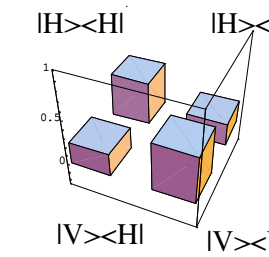
$$\phi_1 = \frac{1}{4} \text{ArcCos} \left[ \sqrt{(2A \cos \phi_1)^2 + 4B^2} \right]$$

$$\phi_2 = \frac{1}{4} \left[ \text{ArcTan} \left[ \frac{2B \cos \phi_1}{2A \cos \phi_1} \right] + \text{ArcTan} \left[ \frac{2B \sin \phi_1}{\sqrt{(2A \cos \phi_1)^2 + 4B^2 \cos^2(\phi_1)}} \right] \right]$$

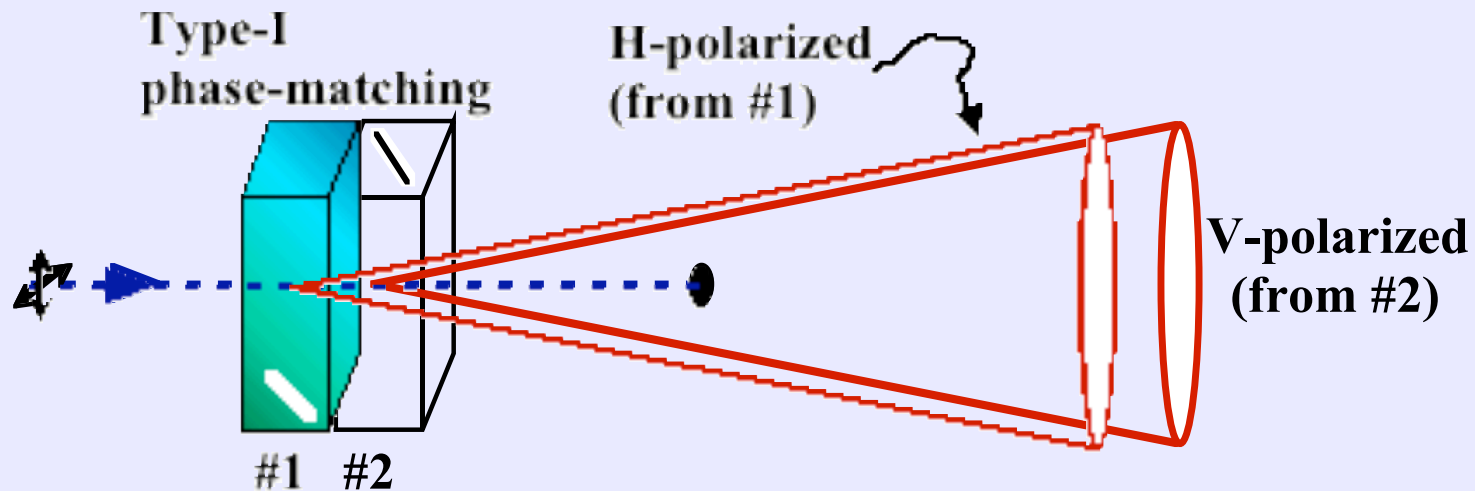
$$\phi_3 = \frac{1}{2} \text{ArcTan} \left[ \frac{2B \cos \phi_1}{2A \cos \phi_1} \right]$$



# Experiment vs. Theory

State	$ H\rangle$	$ \bar{D}\rangle$	Mixed	Partially Mixed
$\square$	$\begin{array}{ c c } \hline 1 & 0 \\ \hline 0 & 0 \\ \hline \end{array}$	$\begin{array}{ c c } \hline \frac{1}{2} & \frac{1}{2} \\ \hline 1 & 1 \\ \hline \end{array}$	$\begin{array}{ c c } \hline \frac{1}{2} & 0 \\ \hline 0 & 1 \\ \hline \end{array}$	$\begin{array}{ c c } \hline 0.5 & 0.25 \\ \hline 0.25 & 0.5 \\ \hline \end{array}$
Theory $\square$				
Experiment $\square$				
Fidelity	0.998	0.996	0.998	0.999

## Two-crystal Source



$$|\psi\rangle = \frac{1}{\sqrt{2}} \left( |H\rangle_1 |H\rangle_2 + e^{i\varphi} |V\rangle_1 |V\rangle_2 \right)$$

**Maximally-entangled state**

## Arbitrary two-qubit pure states

- Via spontaneous parametric downconversion (PDC), our experimental system can generate pure states of the form  $\cos\theta|HH\rangle + \sin\theta e^{i\phi}|VV\rangle$  where the two angles are tunable by waveplate settings before PDC
- We can obtain arbitrary states of the form  $a|HH\rangle + b|HV\rangle + c|VH\rangle + d|VV\rangle$ , by preparing a state with appropriate  $\theta$  and  $\phi$  followed by a series of waveplates in both arms (i.e., local unitary transformations)
- Mathematically, given  $(a,b,c,d)$  (properly normalized), we can find  $(\theta, \phi)$  and local unitaries  $(U^A, U^B)$  such that  $U_1 U_2(\theta, 0, 0, \phi)^T = (a, b, c, d)^T$
- Solution:

$$|\psi\rangle = a|HH\rangle + b|HV\rangle + c|VH\rangle + d|VV\rangle = \sqrt{p}| \tilde{H} \tilde{H} \rangle + \sqrt{1-p}| \tilde{H}^{\perp} \tilde{H}^{\perp} \rangle \quad (\text{Schmidt decomposition})$$

$$\tilde{H} = \sqrt{p}, \tilde{H}^{\perp} = \sqrt{1-p}$$

$$U^A = | \tilde{H} \tilde{H} | + | \tilde{H}^{\perp} \tilde{H}^{\perp} |, U^B = | \tilde{H} \tilde{H} | + | \tilde{H}^{\perp} \tilde{H}^{\perp} |$$

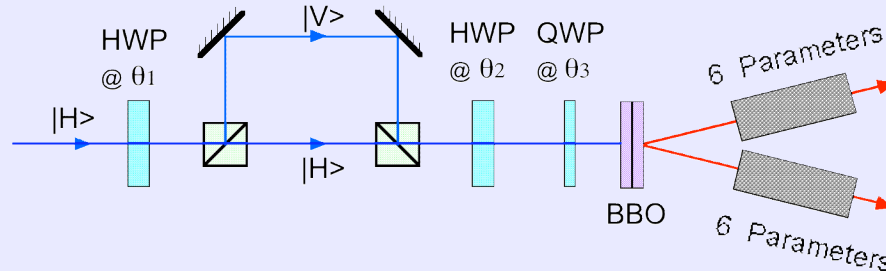
$$|\psi\rangle = U^A U^B (\tilde{H}|HH\rangle + \tilde{H}^{\perp}|VV\rangle)$$

## Arbitrary Two-qubit State Synthesis: Any Mixture/Purity (work in progress)

$$\begin{pmatrix} A_1 & A_4 e^{i A_5} & \dots \\ A_4 e^{-i A_5} & A_2 & \\ & & A_3 \\ & & & 1 - (A_1 + A_2 + A_3) \end{pmatrix}$$

15 free parameters ( $A_1, A_2, \dots$ )

$\Rightarrow$  15 knobs! ( $\theta_1, \theta_2, \dots$ )



$$\theta_1 = f(A_1, A_2, \dots)$$

$\vdots$

**Example:** If we can set each of the parameters to one of 10 values, then we should be able to make  $10^{15}$  states.

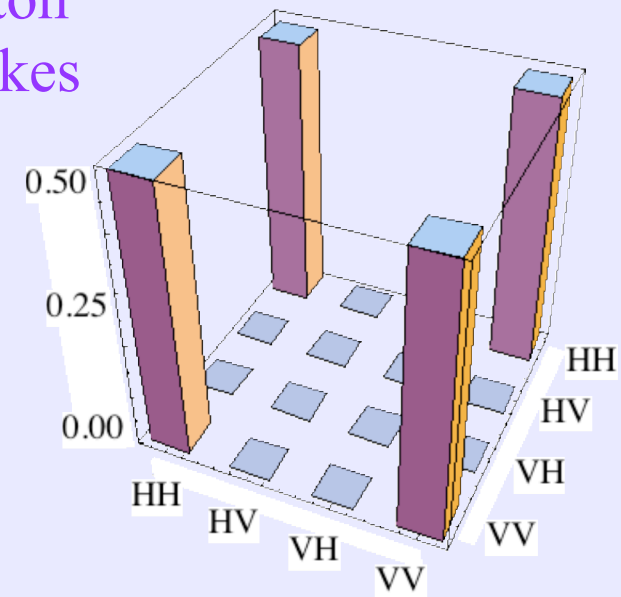
# Quantum State Tomography

Measuring the density matrix for the 2-photon quantum system (c.f., measuring the Stokes parameters for a single photon)

$$\rho_{\text{pure}} = |\psi\rangle\langle\psi|$$

Example:  $|\psi\rangle = \frac{(|HH\rangle + |VV\rangle)}{2}$

$$\rho = \frac{\{|HH\rangle\langle HH| + |VV\rangle\langle VV| + |HH\rangle\langle VV| + |VV\rangle\langle HH|\}}{2}$$



# Experimental Tomography

Single Photon - determine the Stokes parameters

Measure contribution of **H**, **V**, **D**, **R**

Two Photon

Measure polarization correlations:

HH	HV	HD	HR
VH	VV	VD	VR
DH	DV	DD	DR
RH	RV	RD	RR

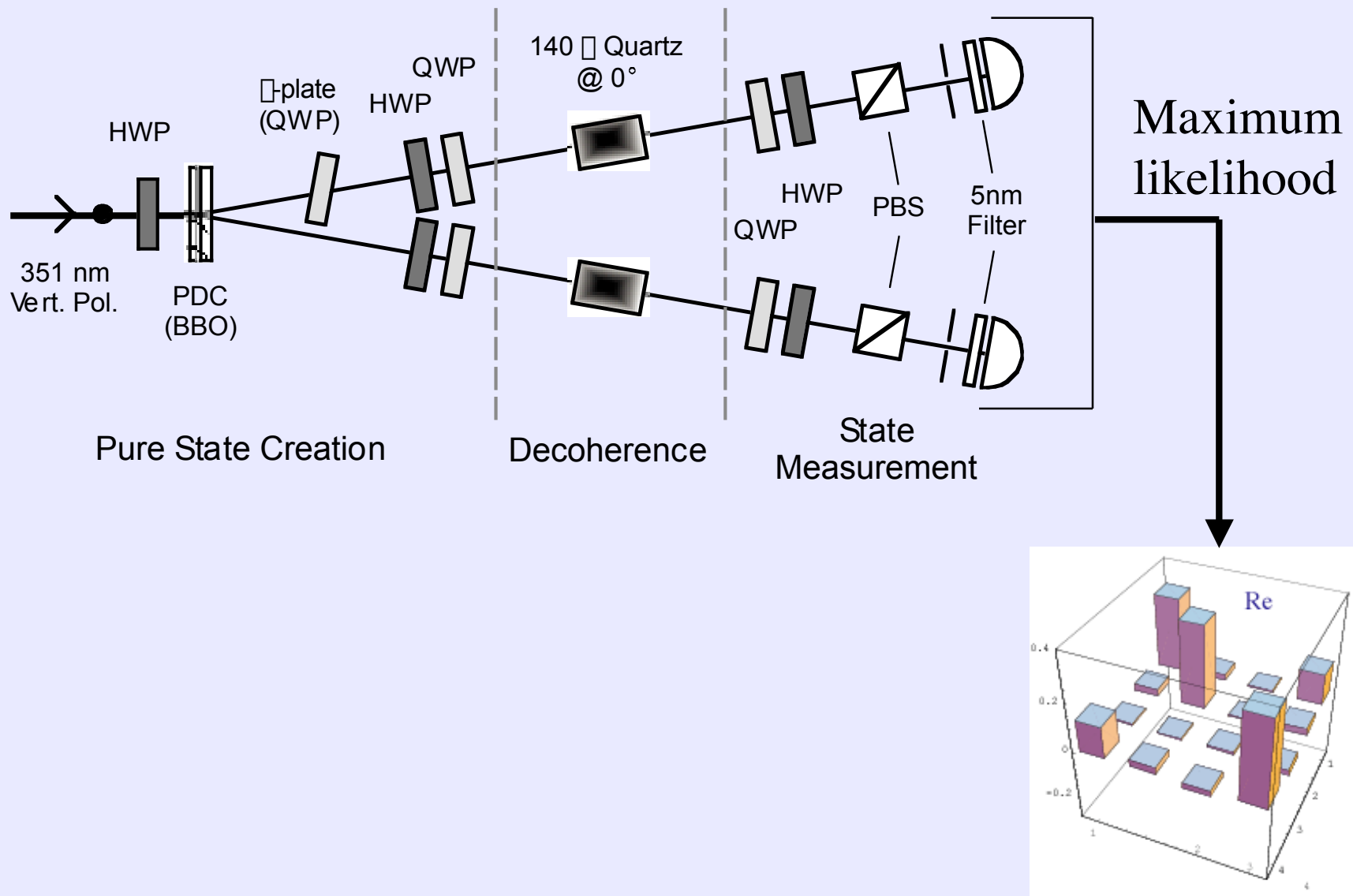
In principle, we can analytically extract the density matrix  $\rho$  from these measurements.

In practice, statistical noise and setting uncertainties will typically lead to a  $\rho$  with *negative* eigenvalues. ☹

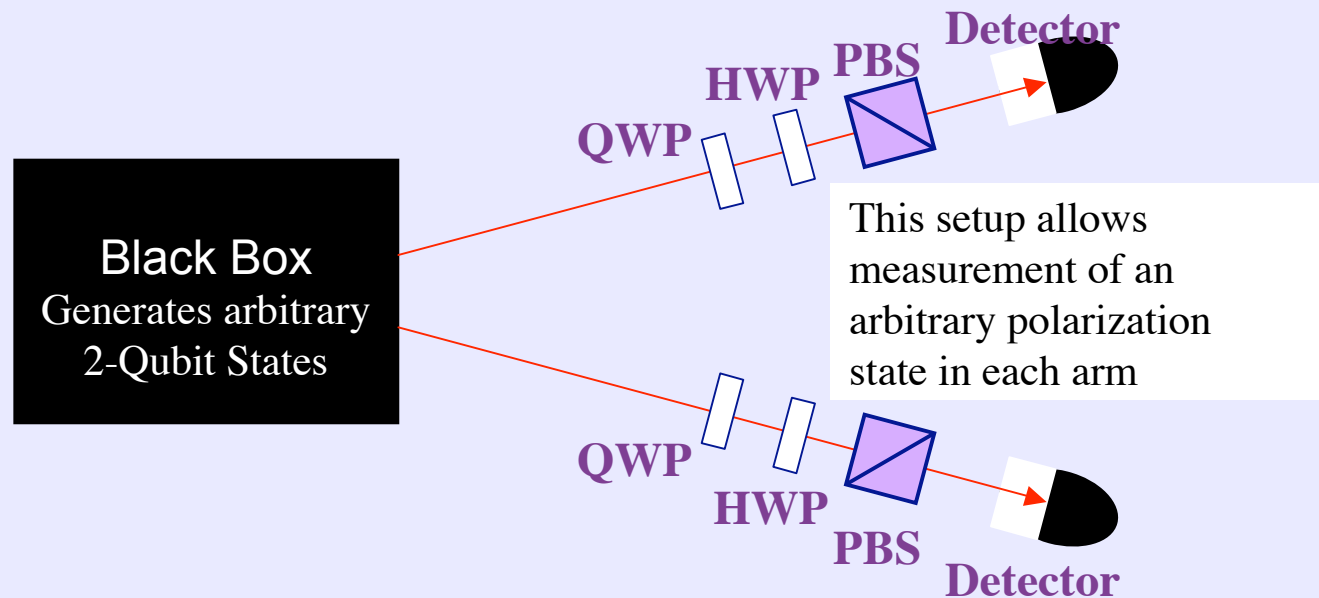
Solution: Use method of maximum likelihood to find the legitimate  $\rho$  most consistent with the measured data.

[D. James et al. [quant-ph/0103121](https://arxiv.org/abs/quant-ph/0103121)]

# 2-qubit state creation & tomography



# Computer Automated Tomography of 2-Qubit Polarization States



Any 2-qubit tomography  
requires 16 of these  
measurements.

## Examples

### Arm 1

H  
H  
D

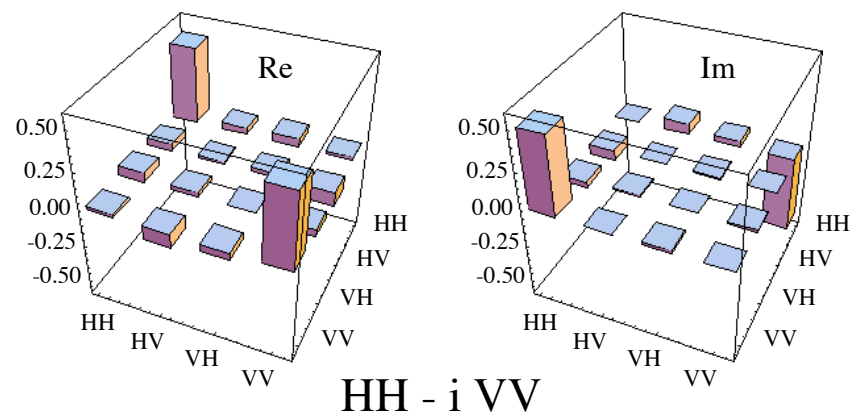
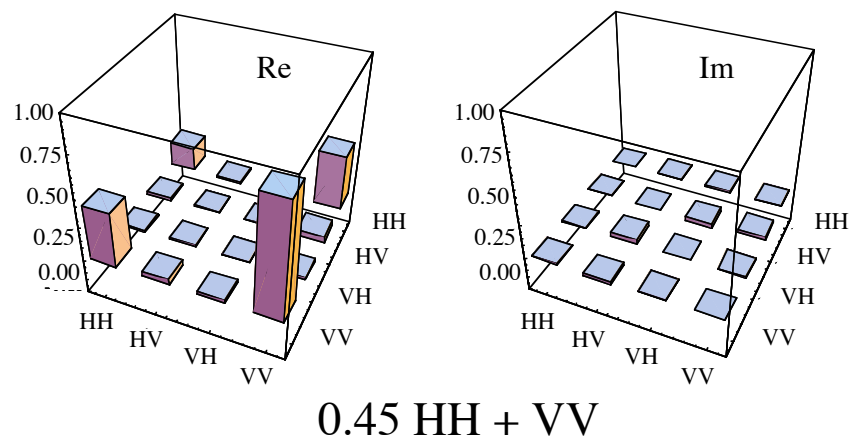
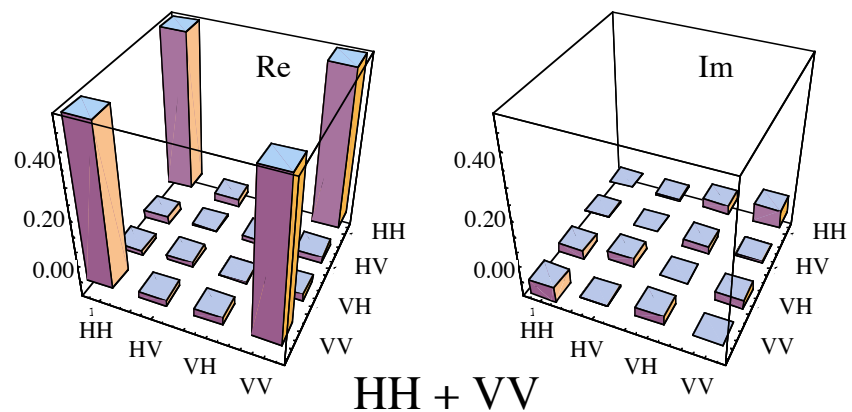
### Arm 2

V  
R  
D

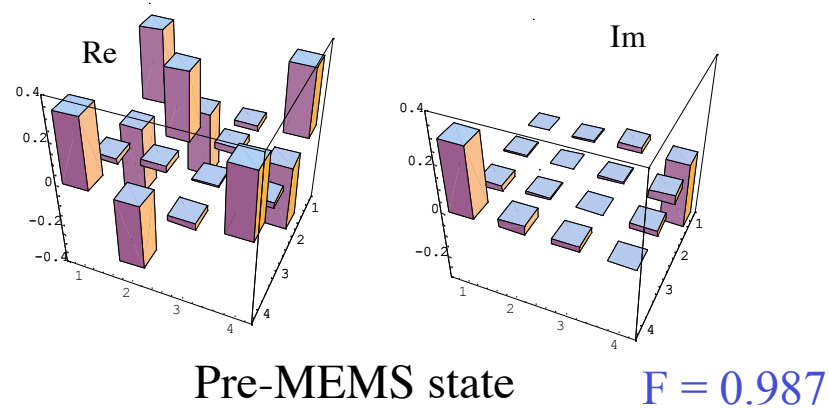
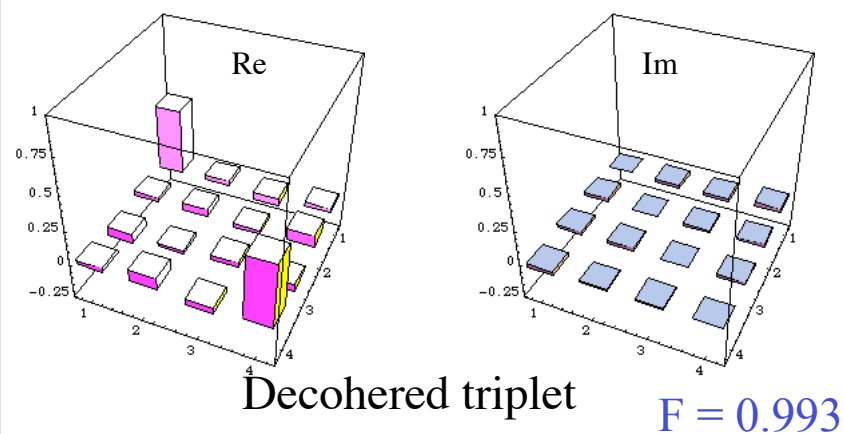
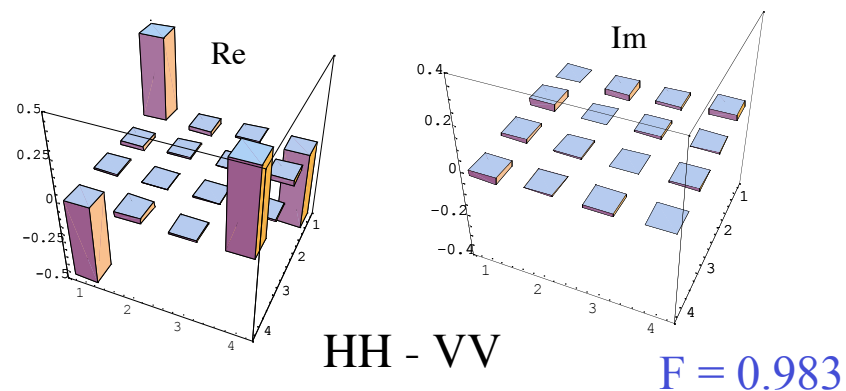
Manual (the old way)	Automated	Automated and Compensated (10 - 100x Brighter)	
		Ultra-Precise	Ultra-Fast
Manually set WP's for each measurement. (100 seconds each)	Computer sets WP's for all measurements. (100 seconds each)	Computer sets WP's for all measurements. (100 seconds each)	Computer sets WP's for all measurements. (10 seconds each)
60 minutes	30 minutes	30 minutes	<b>3 minutes</b>
Counting Error: <b>1.0 %</b>	Counting Error: <b>1.0 %</b>	Counting Error: <b>0.1 – 0.3 %</b>	Counting Error: <b>0.3 – 1.0 %</b>
WP Setting Error: <b>0.5 degrees</b>	WP Setting Error: <b>0.001 degrees</b>	WP Setting Error: <b>0.001 degrees</b>	WP Setting Error: <b>0.001 degrees</b>



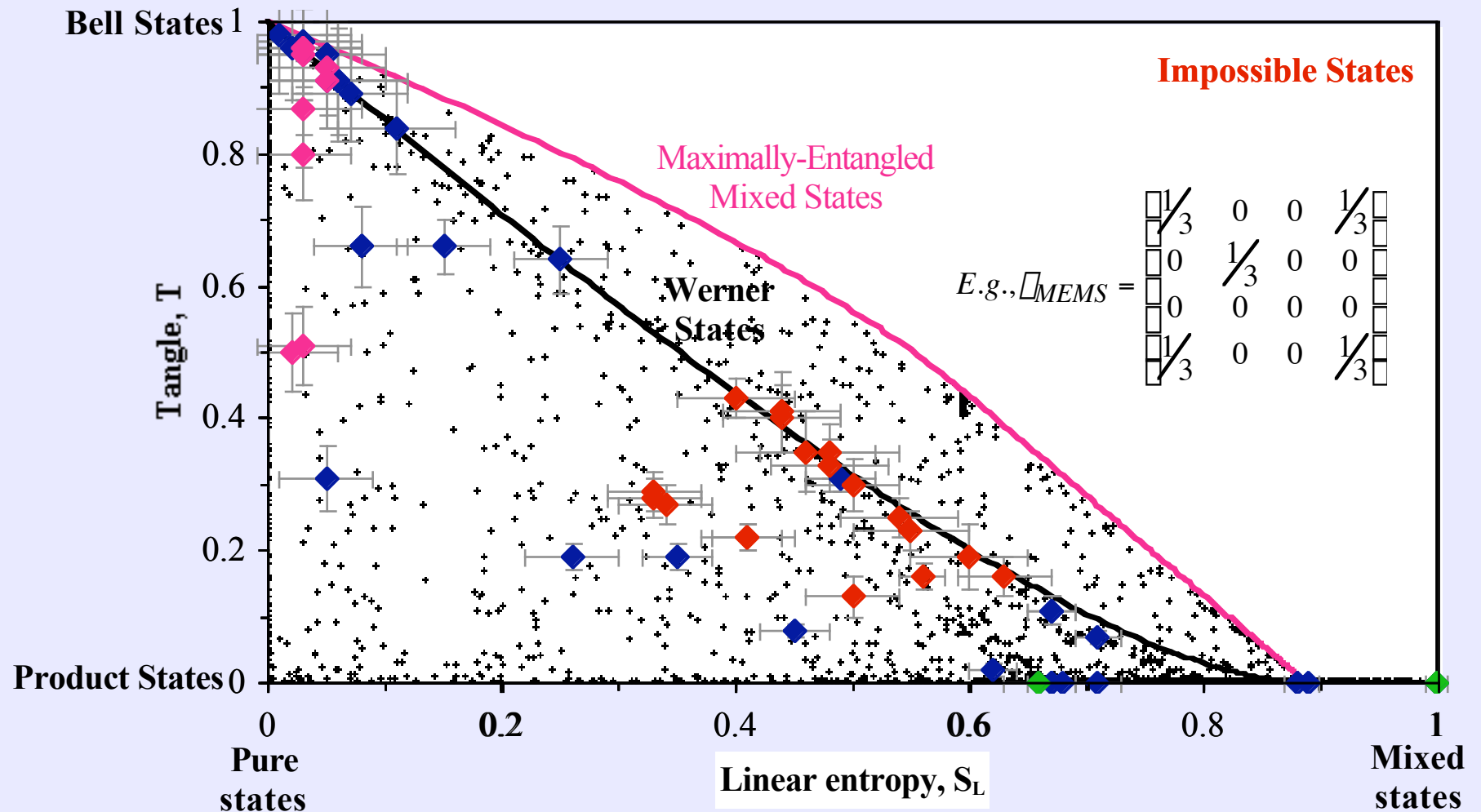
## 2-Photon Quantum Tomography



## From New System:



# 2-parameter characterization of 2-qubit states



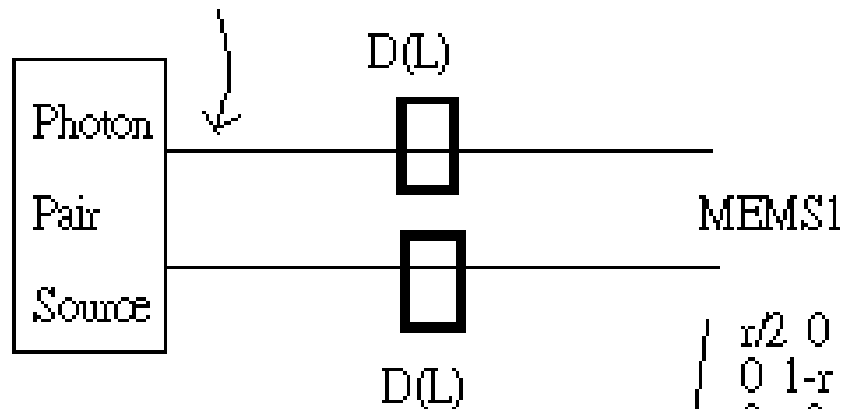
White et al., PRA **65**, 012301 (2001)

Munro et al., PRA **64**, 030302 (2001)

# Making MEMS:

Creation of MEMS 1: one decoherence will do

$$|\psi\rangle = \sqrt{r/2}HH + \sqrt{1-r}HV + \sqrt{r/2}VV$$

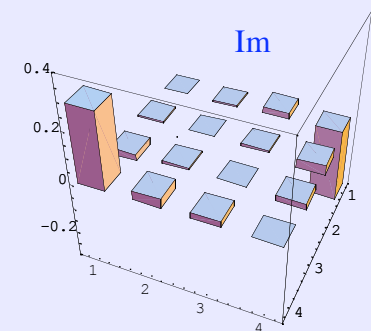
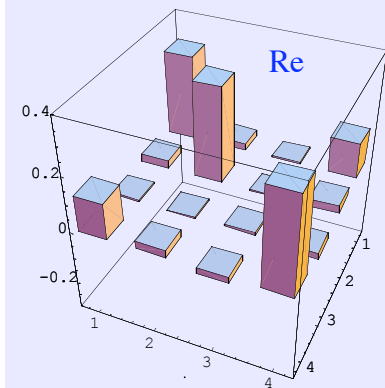


$$\begin{pmatrix} r/2 & 0 & 0 & r/2 \\ 0 & 1-r & 0 & 0 \\ 0 & 0 & 0 & 0 \\ r/2 & 0 & 0 & r/2 \end{pmatrix}$$

$\Delta n L \ll \text{coherence length of pump}$

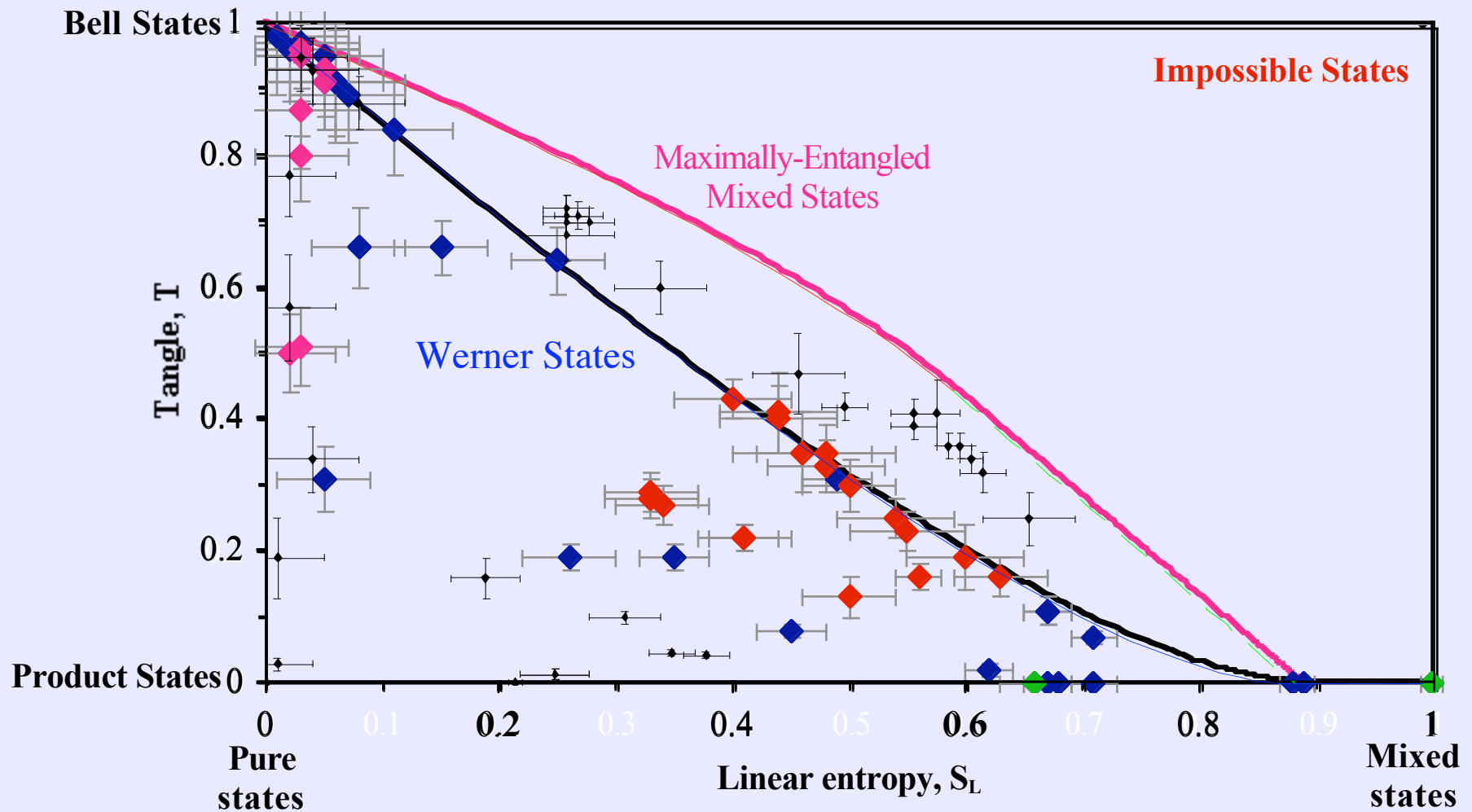
$\Delta n L \gg \text{coherence length of PDC pair}$  e.g.,  $r=2/3$ , get you the  $1/3$  MEMS state

## First MEMS:

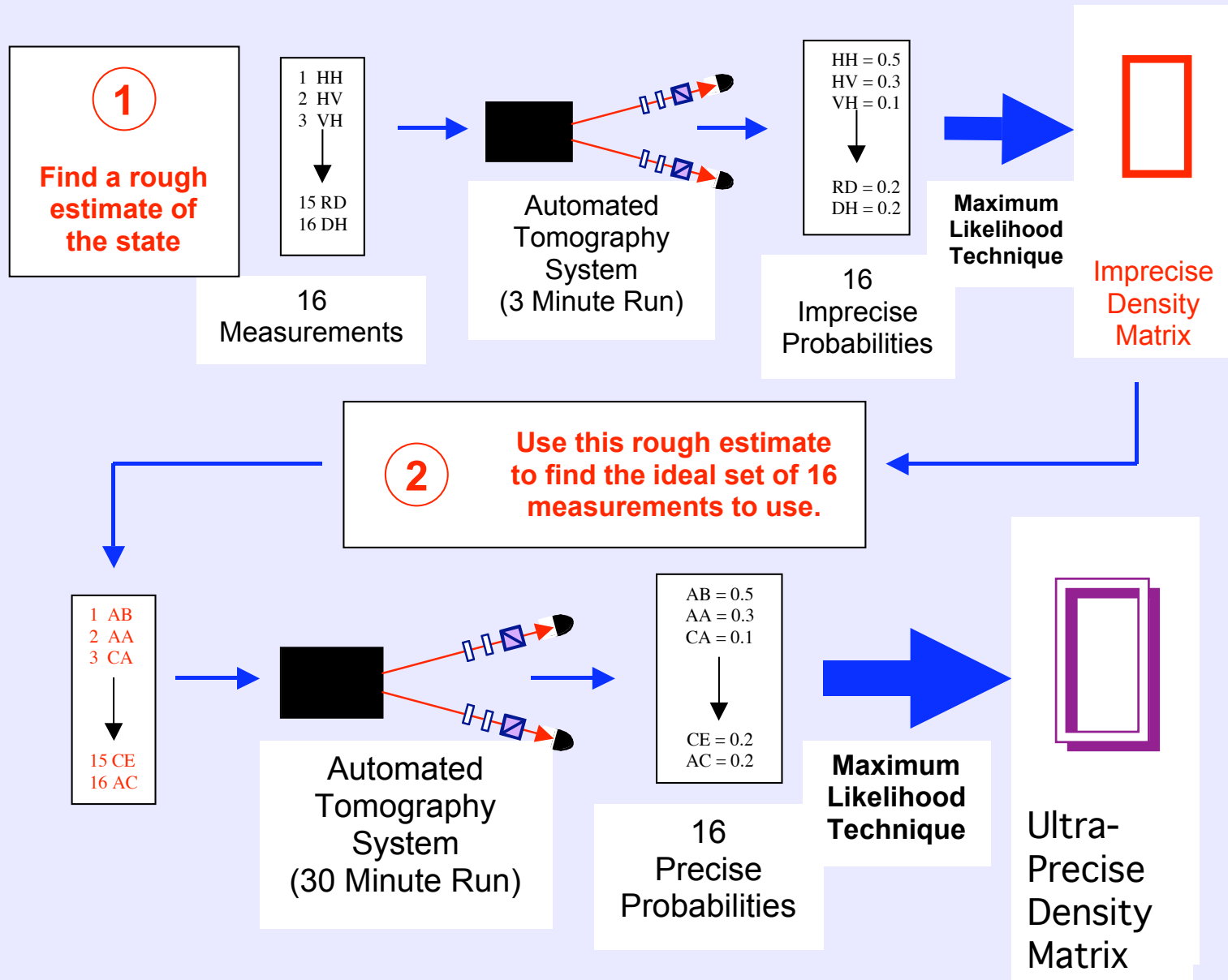


**F = 0.985**

# Beyond Werner States...



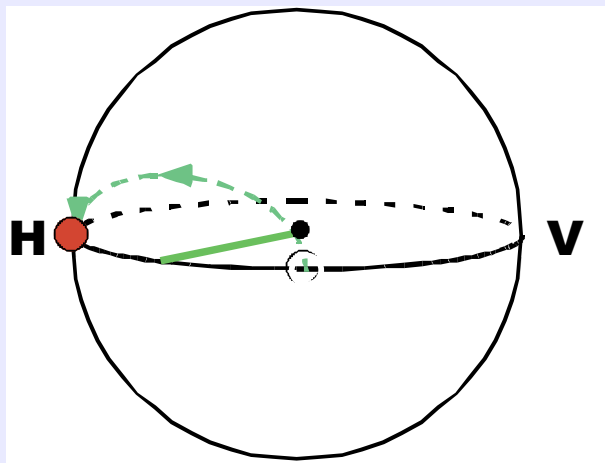
# Adaptive Tomography



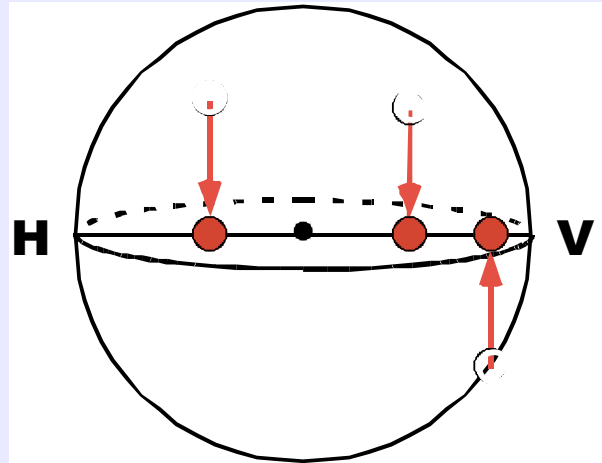
# What is a quantum process?

A completely positive map, which converts an initial quantum state (possibly mixed) into a different quantum state, e.g. unitary transformations and decoherence.

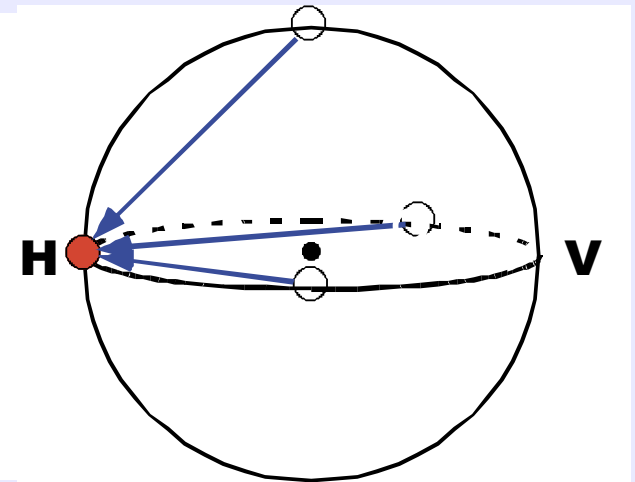
Unitary Transformation:



Decoherence:



Polarizer:

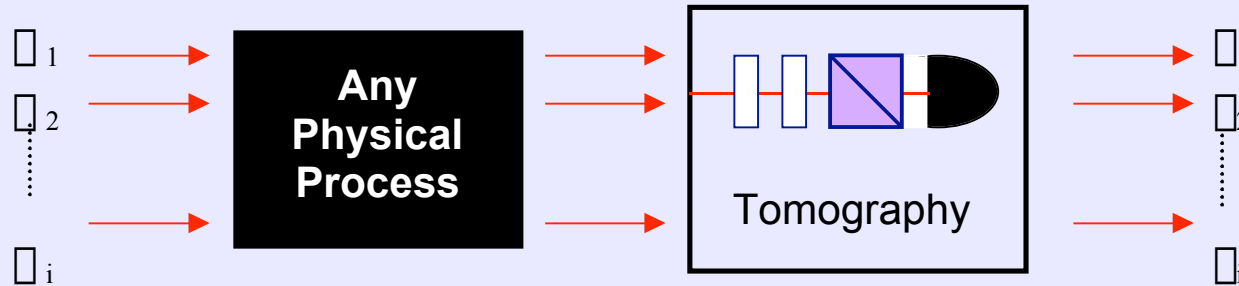


How do we characterize one?

# Quantum Process Tomography

Chuang & Nielsen, JMO 44, 2455 (1997);  
Poyatos, Cirac & Zoller, PRL 78, 390 (1997)

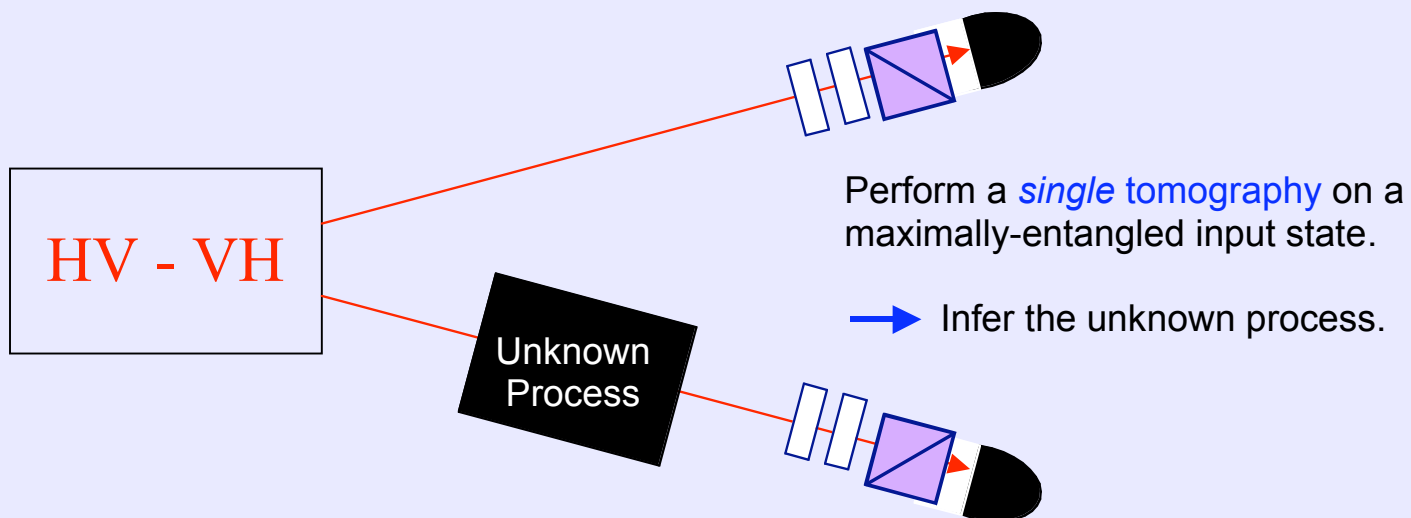
## STANDARD



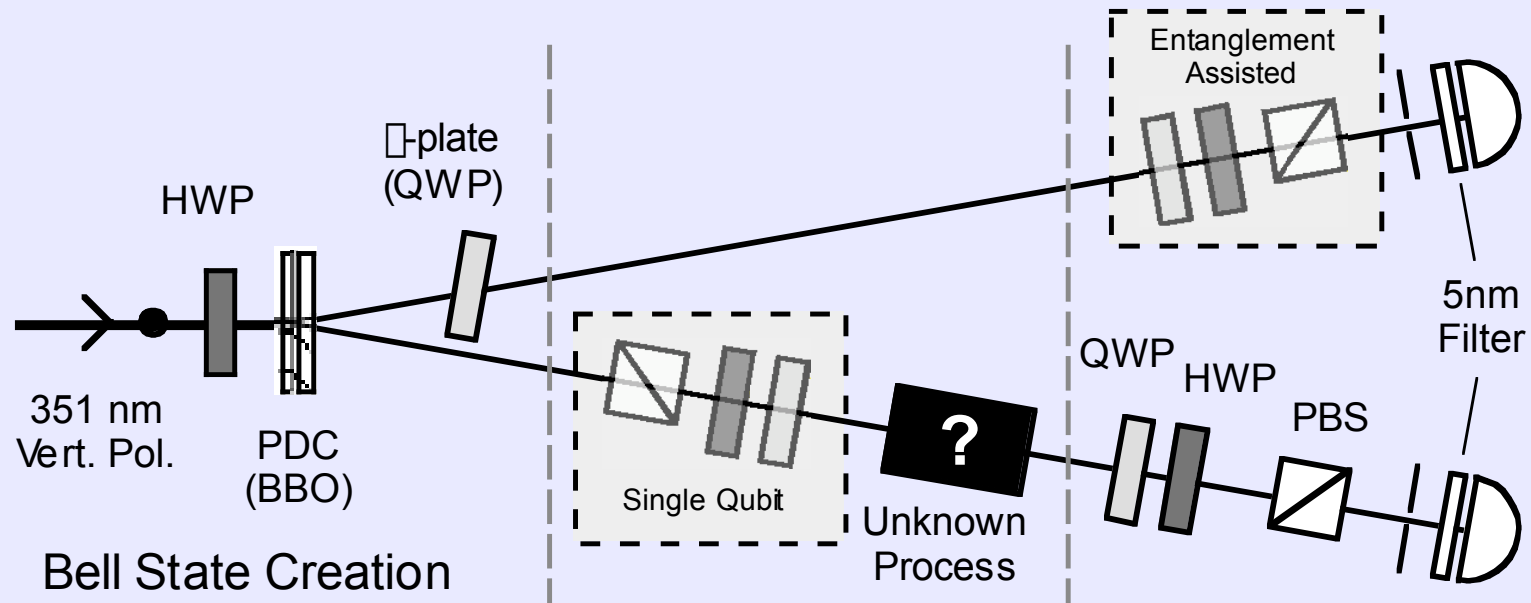
The  $| \square \rangle \rightarrow | \square \rangle$  transforms allow us to reconstruct the unknown process.

## ENTANGLEMENT-ASSISTED

(D'Ariano & Presti; Leung; Nielsen, Thew and White)



# Single-Qubit & Entanglement Assisted Process Tomography





## Single-qubit

## Entanglement-assisted

Process

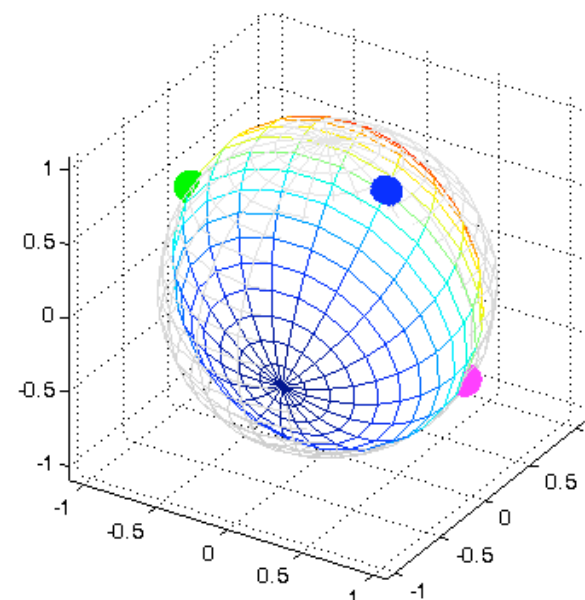
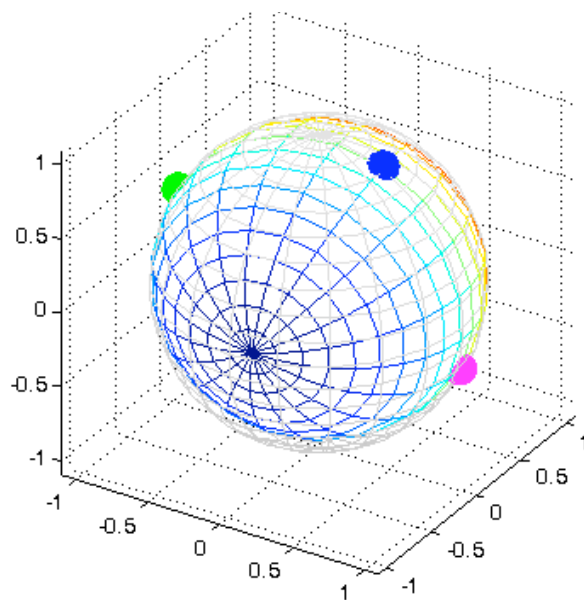
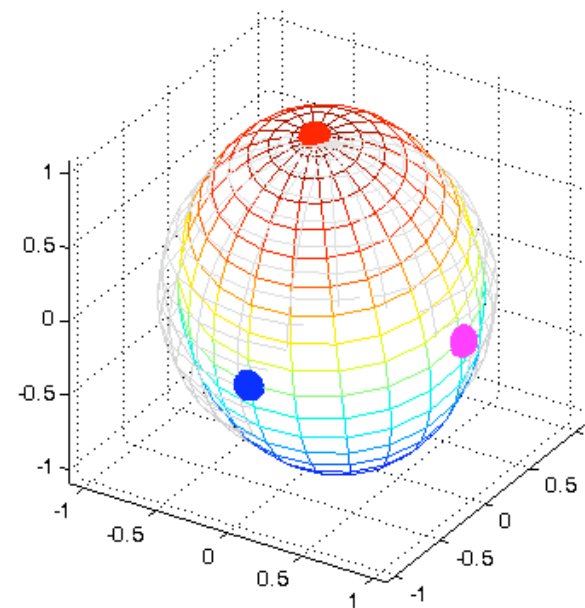
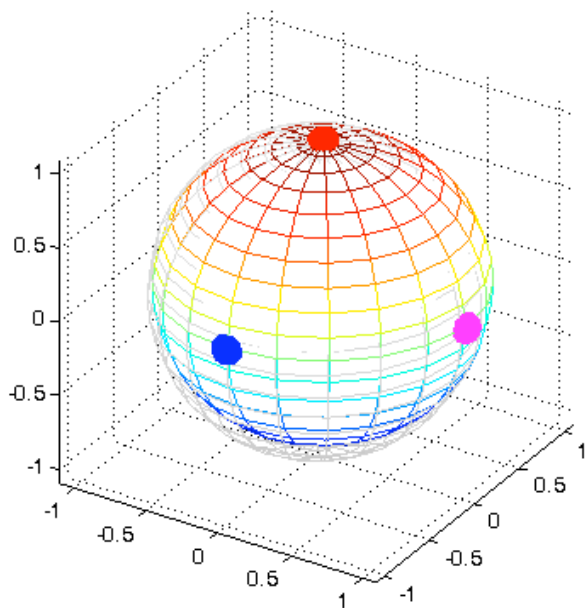
$\langle F \rangle$

Identity:

99.2

Rotate:

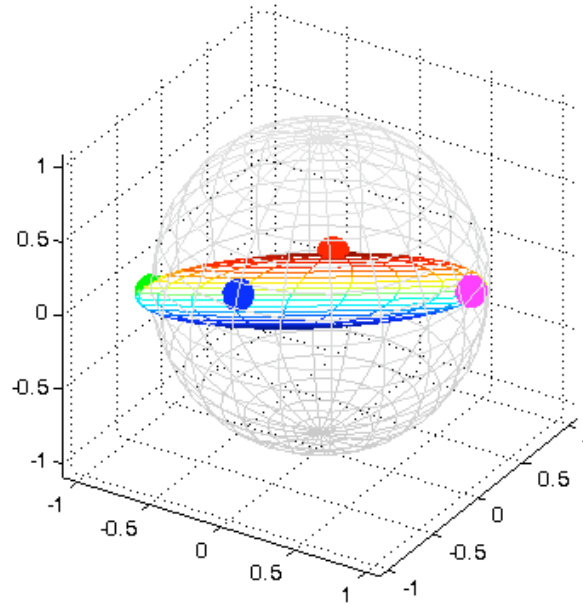
99.1



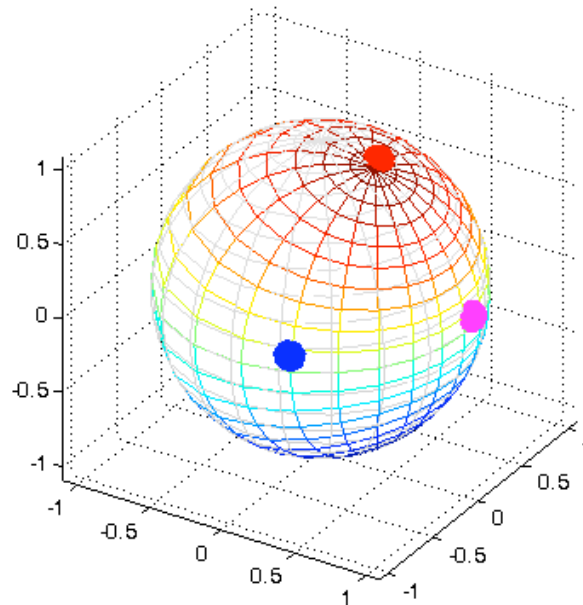
## Single-qubit

Process

Decoherer:



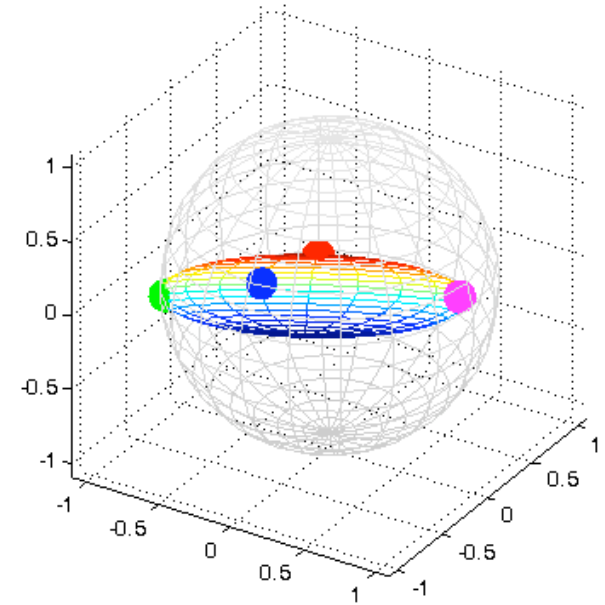
Partial  
polarizer:



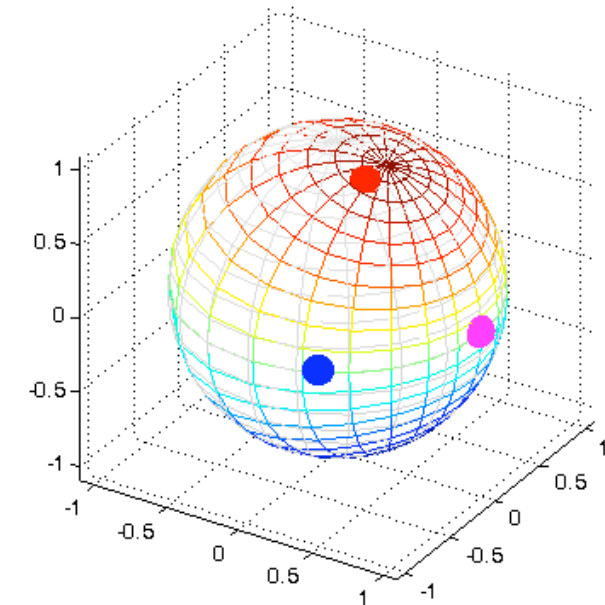
## Entanglement-assisted

$\langle F \rangle$

99.3



98.7

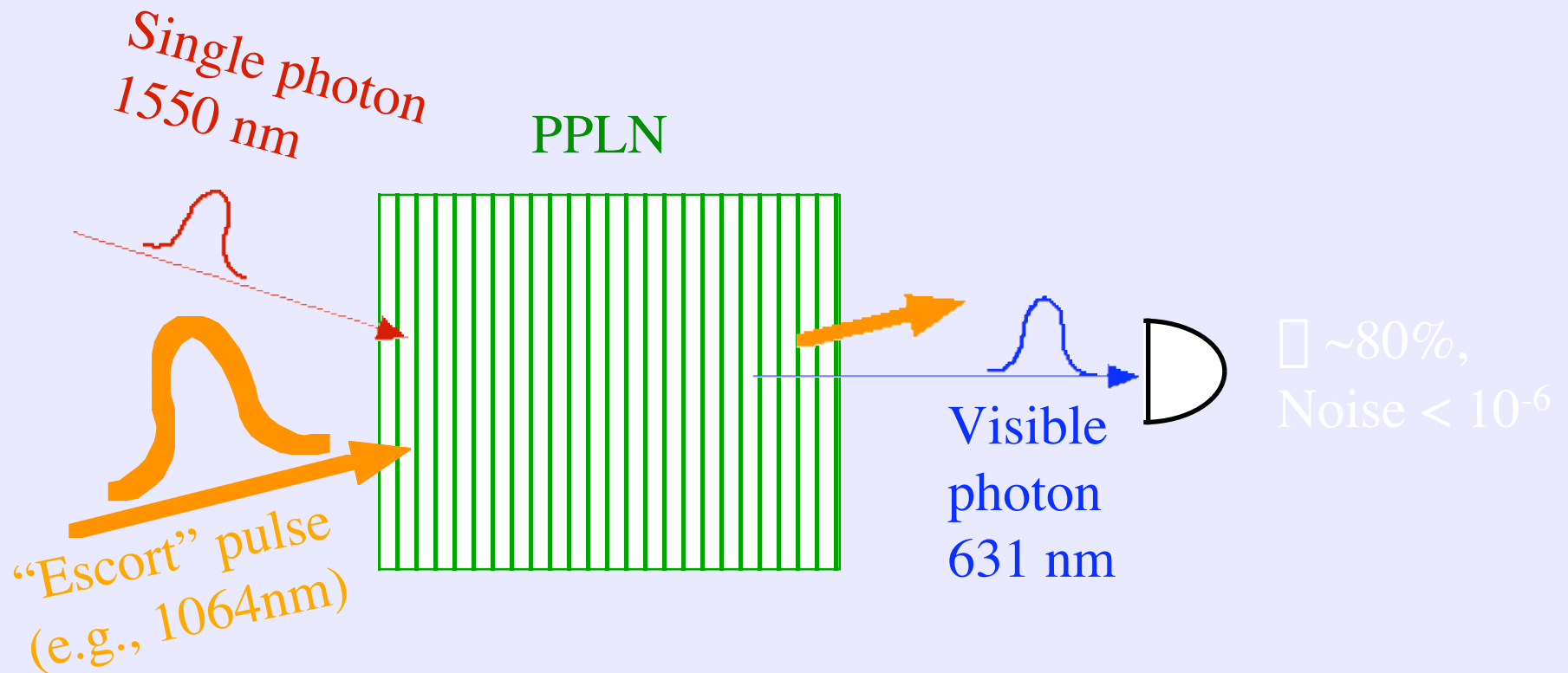


# Quantum Transducer

Convert photon quantum state from one wavelength to another

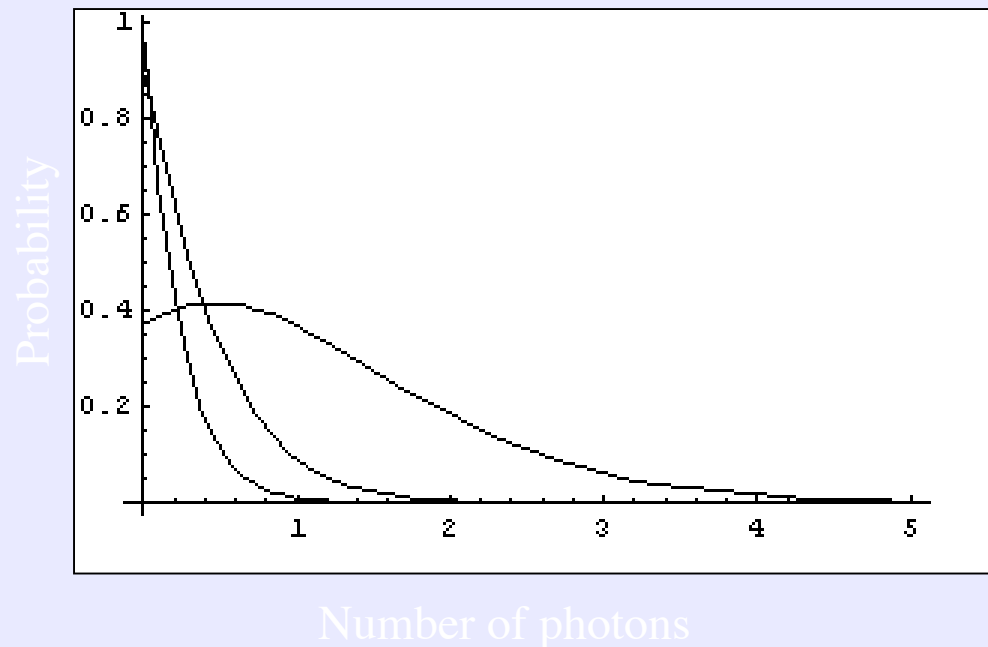
- Capitalize on high-efficiency low-noise single-photon counters (Si APDs) from 500-800nm
- Convert from (telecom) flying qubit to (atomic) storage qubit

Example:



# Deterministic Single Photon Source

Classical (laser) source  $\rightarrow$  Poisson photon number distribution

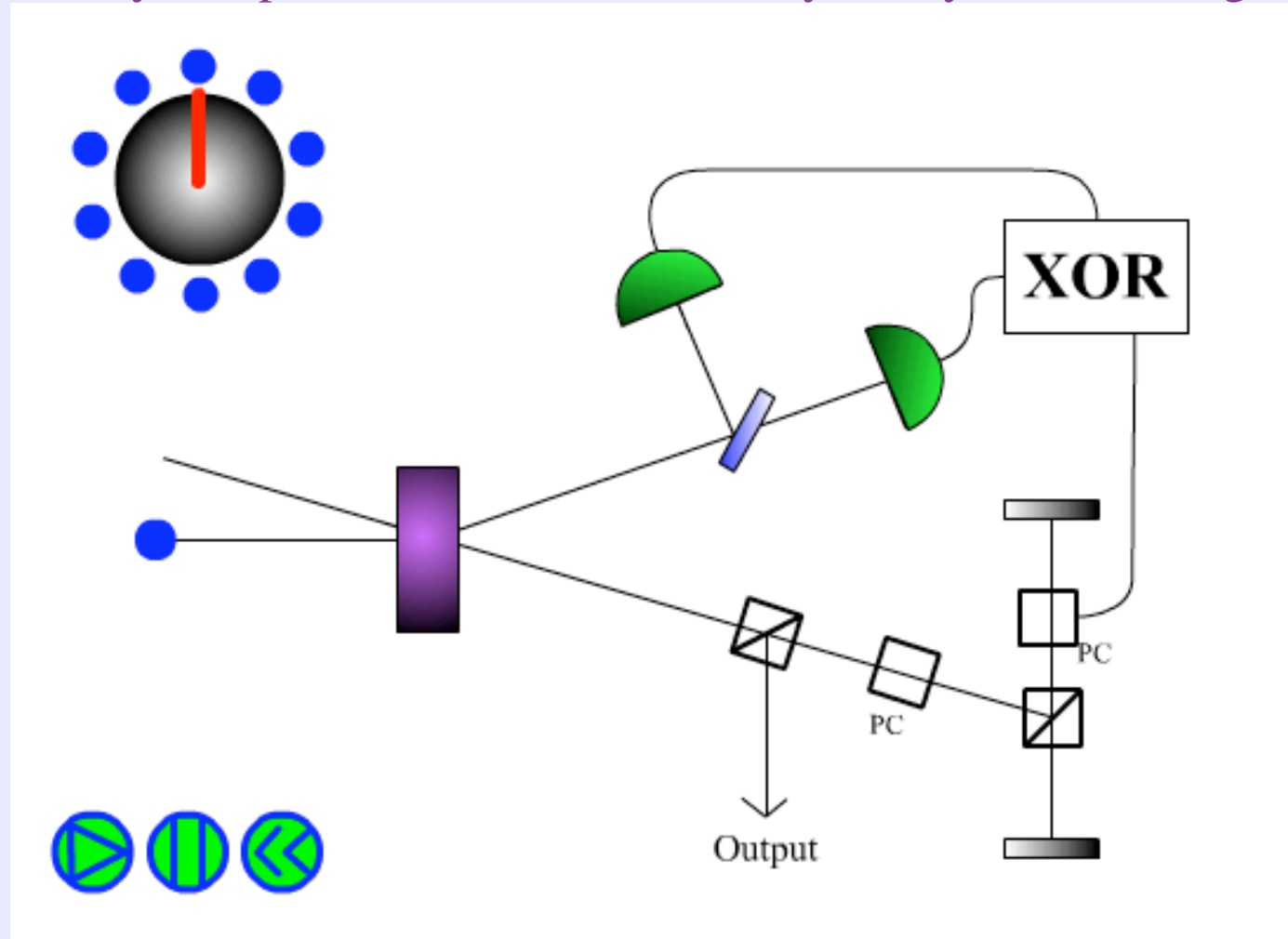


We wish to overcome this, yet retain controllable spectral characteristics

- Optical quantum computing and communication
- Quantum metrology
- Quantum cryptography

# Experimental Design

We collect down-conversion from several pulses, and search for one that has exactly one pair, then store it in a delay cavity until the target time



→ single-photon output at regular intervals  
(~*deterministic* output from *nondeterministic* input!)

# Target efficiencies

## Compare to a classical pulsed laser

37% maximum probability of making one photon

18% probability of making two photons

## Current technology

70% probability of making one photon

3% probability of making two photons

10 kHz rate

## Stage 2 technology (multiple attempts method)

90% probability of making one photon

1.8% probability of making two photons

50 kHz rate

## Future technology (more efficient detectors, lasers, electrooptics)

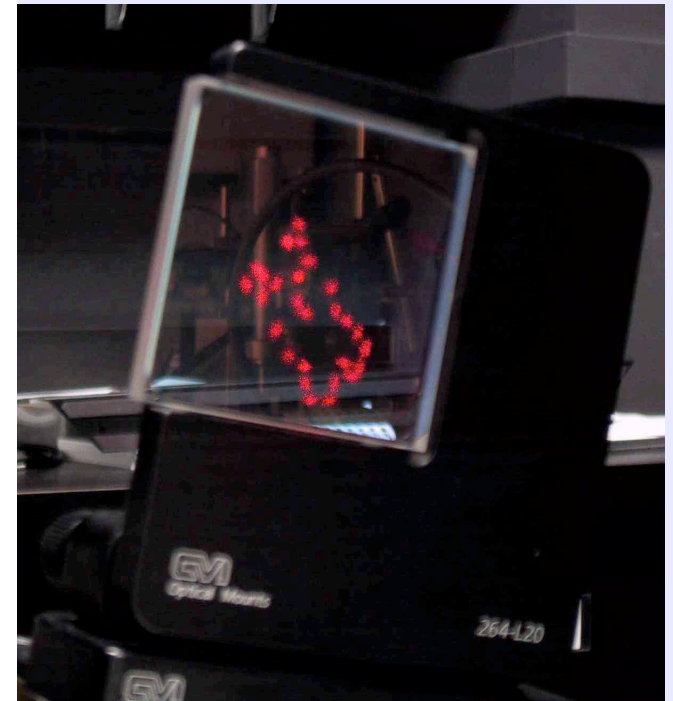
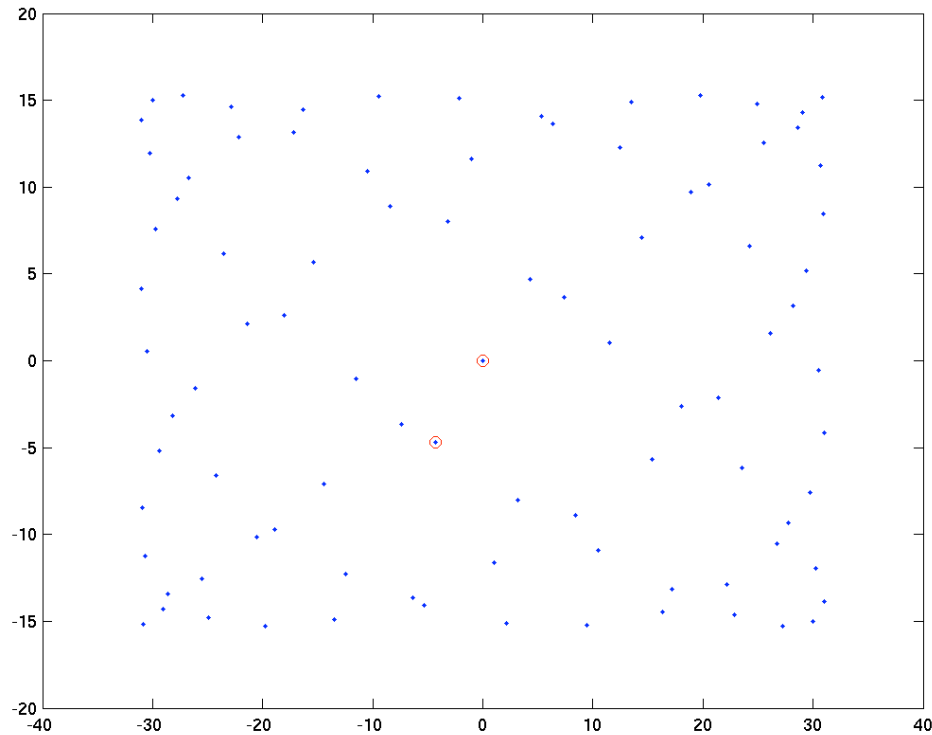
95% probability of making one photon

<1% probability of making two photons

10 MHz rate

# Quantum Memory

In slow systems, the delay line may need to store photons for 1-50  $\mu$ s. Using a specially designed cavity, we can achieve this.



50 pass (300 ns) cavity

This can also be used as a quantum storage in other quantum information applications.

# Summary

- High-fidelity synthesis of arbitrary single-qubit states
- Bright source of tunable 2-qubit states  
Even brighter source to come
- Automated quantum tomography  
Adaptive tomography still to come
- Filling the tangle-entropy plane -- MEMS  
Arbitrary states still to come
- Single qubit/entanglement-assisted process tomography  
Arbitrary process synthesis still to come
- Also in progress:  
Quantum transducer, Single photons, Quantum memory



# Kwiat's *Quantum* Clan

